Waveguides and Resonators

S.M. Riazul Islam, PhD University of Dhaka

Reference Book: Electromagnetics By John D Kraus

- Waveguide: Tx line which can convey EM waves only in higher-order modes.
- Circuit Theory and Field Theory



Can such a pipe convey EM energy

• TE Mode Wave in the infinite Parallel-Plane Tx Line or Guide



- TE: E is everywhere transverse, but H longitudinal , as well as transverse, comp.
- E_v vanishes at the sheets!



Two plane TEM waves traveling in free space in different directions result in maximum field along middle horizontal line (dash double-dot) and zero field along the two dash-dot lines.

- Resultant waves belong to higher order mode.
- TE mode not transmitted unless sufficiently short wavelength: Cutoff Wavelength





$$CB = BD = C'B = \frac{\lambda_0}{4} \qquad CB = \frac{n\lambda_0}{4}$$



 n corresponds to a particular higher order mode. For example, if n=1,

$$\lambda_{oc} = 2b$$
 Cutoff wavelength

$$\theta = \sin^{-1} \frac{\lambda_0}{\lambda_{oc}}$$



$$v_0 = \frac{1}{\sqrt{\mu\epsilon}} \qquad (\text{m s}^{-1})$$

Permeability (Hm⁻¹, Permittivity Fm⁻¹)



- Phase velocity approaches an infinite values as the wavelength is increased towards the cutoff value. And, v approaches v₀.
- Phase velocity of higher order mode wave formed by the sheets is always equal or greater than the velocity in an unbounded medium.
- Energy velocity is always equal or less than the velocity in an unbounded medium.







(a) Infinite-parallel-plane transmission line acting as a waveguide for TE wave. E is in y direction with wave in x direction (out of page). The guide consists of two parallel conducting sheets separated by a distance b. Additional sheets introduced normal to E_y as in (b) results in the hollow rectangular waveguide (c).

- Start w. Maxwell's equation
- Restriction of harmonic variation wrt time
- Harmonic variation and attenuation wrt x
- Select mode TE or TM
- Find four fields in terms of H_x
- Develop Wave equation for H_x
- Solve this Eq. for H_x subject to boundary condtions
- Substitute H_x for other fileds



• Step 1

$$\frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} - \sigma E_x - \epsilon \frac{\partial E_x}{\partial t} = 0 \qquad (1)$$

$$\frac{\partial H_z}{\partial z} - \frac{\partial H_z}{\partial x} - \sigma E_y - \epsilon \frac{\partial E_y}{\partial t} = 0 \qquad (2)$$

$$\frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} - \sigma E_z - \epsilon \frac{\partial E_z}{\partial t} = 0 \qquad (3)$$

$$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} + \mu \frac{\partial H_x}{\partial t} = 0 \qquad (4)$$

$$\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} + \mu \frac{\partial H_y}{\partial t} = 0 \qquad (5)$$

$$\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} + \mu \frac{\partial H_z}{\partial t} = 0 \qquad (6)$$

• Space free from charge

$$\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0$$
(7)
$$\frac{\partial H_x}{\partial x} + \frac{\partial H_y}{\partial y} + \frac{\partial H_z}{\partial z} = 0$$
(8)

• **Step 2-3:** Assume, field comp varies harmonically with both time and distance and also may attenuate with distance

$$E_y = E_1 e^{j\omega t - yx}$$

dH an	
$\frac{\partial H_x}{\partial y} - \frac{\partial H_y}{\partial z} - (\sigma + j\omega\epsilon)E_x = 0$	(10)
$\frac{\partial H_x}{\partial z} + \gamma H_z - (\sigma + j\omega\epsilon)E_y = 0$	(11)
$-\gamma H_y - \frac{\partial H_x}{\partial y} - (\sigma + j\omega\epsilon)E_z = 0$	(12)
$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} + j\omega\mu H_x = 0$	2 (13)
$\frac{\partial E_x}{\partial z} + \gamma E_z + j\omega\mu H_y = 0$	(14)
$-\gamma E_y - \frac{\partial E_x}{\partial y} + j\omega\mu H_z = 0$	(15)
$-\gamma E_x + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0$	(16)
$-\gamma H_x + \frac{\partial H_y}{\partial y} + \frac{\partial H_z}{\partial z} = 0$	(17)

$Z=-j\omega\mu$	(Ωm^{-1})
$Y = \sigma + j\omega\epsilon$	(Ö m ⁻¹)

$\frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} - YE_x = 0$	(20)
$\frac{\partial H_x}{\partial z} + \gamma H_z - Y E_y = 0$	(21)
$-\gamma H_y - \frac{\partial H_x}{\partial y} - YE_z = 0$	(22)
$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} - ZH_x = 0$. (23)
$\frac{\partial E_x}{\partial z} + \gamma E_z - ZH_y = 0$	(24)
$-\gamma E_y - \frac{\partial E_x}{\partial y} - ZH_z = 0$	(25)
$-\gamma E_x + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0$	(26)
$-\gamma H_x + \frac{\partial H_y}{\partial y} + \frac{\partial H_z}{\partial z} = 0$	(27)

• Step 4: TE Mode

∂H₂ ∂H,	
$\frac{\partial y}{\partial z} = 0$	(28)
$\frac{\partial H_x}{\partial z} + \gamma H_z - Y E_y = 0$	(29) _L
$-\gamma H_y - \frac{\partial H_x}{\partial y} - YE_z = 0$	(30)
$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} - ZH_x = 0$	(31)
$\gamma E_z - ZH_y = 0$	(32)~
$-\gamma E_y - ZH_z = 0$	(33) ~
$\frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0$	(34)
$-\gamma H_x + \frac{\partial H_y}{\partial y} + \frac{\partial H_z}{\partial z} = 0$	(35)

• Step 5:

$$\frac{E_z}{H_y} = -\frac{E_y}{H_z} = \frac{Z}{\gamma} \qquad (\Omega) \qquad (3)$$

Transverse-wave impedance

$$Z_{yz} = \frac{E_y}{H_z} = -\frac{E_z}{H_y} = -\frac{Z}{\gamma} = \frac{j\omega\mu}{\gamma} \qquad (\Omega)$$

$$H_{y} = \frac{-1}{\gamma - YZ_{yz}} \frac{\partial H_{x}}{\partial y}$$
(38)

$$H_{z} = \frac{-1}{\gamma - YZ_{yz}} \frac{\partial H_{x}}{\partial z}$$
(39)

Now, substituting (39) into (37), we obtain

$$E_{y} = \frac{-Z_{yz}}{\gamma - YZ_{yz}} \frac{\partial H_{x}}{\partial z}$$

and substituting (38) into (37) gives

$$E_z = \frac{Z_{yz}}{\gamma - YZ_{yz}} \frac{\partial H_x}{\partial y} \qquad (41)$$

(40)

Step 6: 35:<38, 39 $-\gamma H_{x} - \frac{1}{\gamma - YZ_{yz}} \left(\frac{\partial^{2} H_{x}}{\partial y^{2}} + \frac{\partial^{2} H_{x}}{\partial z^{2}} \right) = 0 \quad (42)$ or $\frac{\partial^{2} H_{x}}{\partial y^{2}} + \frac{\partial^{2} H_{x}}{\partial z^{2}} + \gamma(\gamma - YZ_{yz})H_{x} = 0 \quad (43)$ Putting $k^{2} = \gamma(\gamma - YZ_{yz})$ reduces (43) to $\frac{\partial^{2} H_{x}}{\partial y^{2}} + \frac{\partial^{2} H_{x}}{\partial z^{2}} + k^{2} H_{x} = 0 \quad (44)$

(45)

(46)

(47)

• Step 7: where Y = a function of y only, that is, Y = f(y)Z = a function of z only[†]



Substituting (45) in (44) gives

$$Z\frac{d^2Y}{dy^2} + Y\frac{d^2Z}{dz^2} + k^2YZ = 0$$

Dividing by YZ to separate variables gives

$$\frac{1}{Y}\frac{d^2Y}{dy^2} + \frac{1}{Z}\frac{d^2Z}{dz^2} = -k^2$$

of u at

$$\frac{1}{Y}\frac{d^{2}Y}{dy^{2}} = -A_{1}$$
(48)
$$\frac{1}{Z}\frac{d^{2}Z}{dz^{2}} = -A_{2}$$
(49)

$$Y = c_1 \sin b_1 y$$

Substituting (51) in (48) yields

$$b_1 = \sqrt{A_1} \tag{52}$$

Hence (51) is a solution provided (52) is fulfilled. Another solution is

$$Y = c_2 \cos b_1 y \tag{53}$$

f (51) and (53) are each a solution for Y, their sum is also a solution, or

$$Y = c_1 \sin \sqrt{A_1} y + c_2 \cos \sqrt{A_1} y$$
 (54)



$$H_x(y, z) = H_0 \cos \frac{n\pi y}{y_1} \cos \frac{m\pi z}{z_1}$$
 (59)

Multiplication by constant factor:

$$H_x(y, z, x, t) = H_0 \cos \frac{n\pi y}{y_1} \cos \frac{m\pi z}{z_1} e^{-\gamma x}$$

• Step 8:

$$H_{y} = \frac{\gamma H_{0}}{k^{2}} \frac{n\pi}{y_{1}} \sin \frac{n\pi y}{y_{1}} \cos \frac{m\pi z}{z_{1}} e^{-\gamma x}$$
(61)

$$H_{z} = \frac{\gamma H_{0}}{k^{2}} \frac{m\pi}{z_{1}} \cos \frac{n\pi y}{y_{1}} \sin \frac{m\pi z}{z_{1}} e^{-\gamma x}$$
(62)

$$E_{y} = \frac{\gamma Z_{yz} H_{0}}{k^{2}} \frac{m\pi}{z_{1}} \cos \frac{n\pi y}{y_{1}} \sin \frac{m\pi z}{z_{1}} e^{-\gamma x}$$
(63)

$$E_{z} = -\frac{\gamma Z_{yz} H_{0}}{k^{2}} \frac{n\pi}{y_{1}} \sin \frac{n\pi y}{y_{1}} \cos \frac{m\pi z}{z_{1}} e^{-\gamma x}$$
(64)

If m=1, n=0: H_x, H_z, E_y

- For example: E_y has a sinusoidal variation across the guide (z-direction), being a max in the center and zero at the walls, and has no variation as a function of y.
- m, n → number of half-cycle variation wrt z and y respectively!

• TE₁₀

$$E_{x} = 0 ext{ TE mode requirement}$$

$$E_{y} = \frac{\gamma Z_{yz} H_{0}}{k^{2}} \frac{\pi}{z_{1}} \sin \frac{\pi z}{z_{1}} e^{-\gamma x}$$

$$E_{z} = 0$$

$$H_{x} = H_{0} \cos \frac{\pi z}{z_{1}} e^{-\gamma x}$$

$$H_{y} = 0$$

$$H_{z} = \frac{\gamma H_{0}}{k^{2}} \frac{\pi}{z_{1}} \sin \frac{\pi z}{z_{1}} e^{-\gamma x}$$
(65)



• General Significance:

$$\left(\frac{n\pi}{y_1}\right)^2 + \left(\frac{m\pi}{z_1}\right)^2 = k^2 \tag{67}$$

 $k^2 = \gamma^2 - j\omega\mu(\sigma + j\omega\epsilon) \tag{68}$

For Lossless dielectric medium:

$$\gamma = \sqrt{\left(\frac{n\pi}{y_1}\right)^2 + \left(\frac{m\pi}{z_1}\right)^2 - \omega^2 \mu \epsilon}$$
(69)

At low frequencies, ω small, γ real, guide opaque, wave does not propagate
 At an intermediate frequency, ω intermediate, γ = 0, transition condition (cutoff)
 At high frequencies, ω large, γ imaginary, guide transparent, wave propagates